

# Absorption and Scattering

Definitions – Sometimes it is not clear which process is taking place

# Absorption and Scattering

- Scattering: Photon interacts with the scatterer and emerges in a new direction with (perhaps) a slightly different frequency
  - No destruction of the photon in the sense that energy is added to the kinetic pool
- Absorption: Photon interacts and is destroyed and its energy is converted (at least partially) into kinetic energy of the particles composing the gas.
- Scattering is mostly independent of the thermal properties of the gas and depends mostly on the radiation field.

# Absorption and Scattering

- Absorption processes depend on the thermodynamic properties of the gas as they feed energy directly into the gas.
- NB: thermal emission couples the thermodynamic quantities directly to the radiation field!

# Scattering Processes

- Bound-Bound Transitions followed by a direct reverse transition:
  - $a \rightarrow b : b \rightarrow a$
  - There can be a small  $\Phi E$  as the energy states have finite width (substates)
  - $\Phi E \sim 0$  and “only” the direction changed!
- Photon scattering by a free electron (Thomson) or molecule (Rayleigh)
  - Thomson:  $0.6655 (10^{-24}) N_e$
  - Rayleigh:  $\odot \lambda^{-4}$

# Absorption Processes

- Photoionization / bound-free transition
  - Excess energy goes into KE
  - Inverse Process: Radiative Recombination
- Free-Free Absorption: an electron moving in the field of an ion absorbs a photon causing a shift in the “orbit” (hyperbolic).
  - Inverse Process: Bremsstrahlung
- Bound-Bound Photoexcitation followed immediately by a collisional de-excitation  $\implies$  photon energy shared by partners and goes into the kinetic pool.
  - Inverse Process: -----
- Photo-excitation followed immediately by a collisional ionization.
  - Inverse Process: Collisional Recombination

# Caveats

- These lists are not exhaustive. In many cases the line between the two processes is not clear, especially with bound-free events.

# Absorption Coefficient

- Absorption Coefficient  $\kappa_\nu$  per gram of stellar material such that:
  - A differential element of material of cross section  $dA$  and length  $ds$  absorbs an amount of energy from a beam of specific intensity  $I_\nu$  (incident normal to the ends of the element):
  - $dE_\nu = \rho \kappa_\nu I_\nu d\omega dt dA ds dv$

# Scattering Coefficient

- Similarly,  $\sigma_v$  governs the amount of energy scattered out of the beam.

$$dE_v = \rho \sigma_v I_v d\omega dt dA ds dv$$

- We assume that both  $\kappa_v$  and  $\sigma_v$  have no azimuthal dependence. The combined effect of  $\kappa_v$  and  $\sigma_v$  is to remove energy from the beam.
- The total extinction coefficient is:  $K_v = \kappa_v + \sigma_v$
- $K_v$  is called the “mass absorption coefficient”



# The Simple Slab

$$dI_\nu = -\rho\kappa_\nu I_\nu dz$$

$$I_\nu: \text{ergs Hz}^{-1} \text{ s}^{-1} \text{ cm}^{-1} \text{ sterad}^{-1}$$

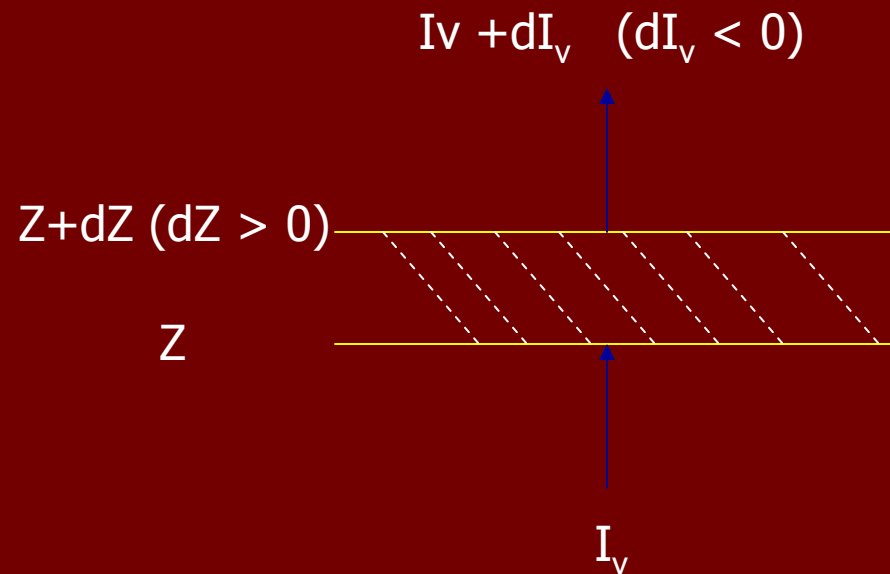
$$\kappa_\nu: \text{cm}^2 \text{ gm}^{-1}$$

$$\rho\kappa_\nu : \text{cm}^{-1}$$

$$\frac{dI_\nu}{I_\nu} = -\kappa_\nu \rho dz$$

$$\ln I_\nu \Big|_{I_{\nu_1}}^{I_{\nu_2}} = -\int_{z_1}^{z_2} \kappa_\nu \rho dz$$

$$I_{\nu_2} = I_{\nu_1} \exp\left(-\int_{z_1}^{z_2} \kappa_\nu \rho dz\right)$$



# The Simple Slab Continued

- In the above 1 and 2 are arbitrary limits
- Let us now integrate over the slab (assign physically meaningful values to the integration limits). We want the “output” at the top ( $Z=0$ ) coming from depth  $X$  in the slab (star?!)
  - $I_{v1}$  = Initial intensity =  $I_v(X)$
  - $I_{v2}$  = Output Intensity
  - $Z_1$  = Level of Origin
  - $Z_2$  = Level of Output (= 0)

$$I_v(0) = I_v(X) \exp\left(-\int_x^0 \kappa_v \rho dz\right)$$

# The Optical Depth

The Dimensionless Quantity  $\tau$

$$d\tau_\nu = -\kappa_\nu \rho dx$$

$$\tau_{\nu_2} - \tau_{\nu_1} = -\int_{x_1}^{x_2} \kappa_\nu \rho dx$$

Note that  $\tau$  and  $X$  increase in opposite directions. Boundary Conditions: At the surface  $\tau_{\nu_1} \equiv 0$  and at the “base”  $X_1 \equiv 0$ . Therefore:

$$\tau_{\nu_2} = -\int_0^{x_2} \kappa_\nu \rho dx$$

$$-\tau_\nu = -\int_x^0 \kappa_\nu \rho dx$$

$\therefore$

$$I_\nu(0) = I_\nu(x)e^{-\tau_\nu}$$

# Optical Depth

## Surfaces and Probability

- The optical depth of a slab determines how much light escapes from a given level:
  - If  $\tau = 2$  then  $I = I_0 e^{-2} \implies 0.135$  of  $I_0$  escapes
- The usual definition of “continuum” (AKA the “surface” of a star) is where a photon has a 50% chance of escaping:
  - $e^{-\tau} = 0.5 \implies \tau = 0.693$
- An exercise in atmospheric extinction. A cloud can easily contribute 3 magnitudes of extinction:
  - $1/(2.512)^3 = e^{-\tau} \implies \tau = 2.763$ .

# Emission Processes

If there is a sink then there must be a source!

- Emission coefficient:  $j_\nu$ 
  - $dE_\nu = \rho j_\nu d\omega dv dt dA ds$
  - Or per unit mass in an incremental volume:
  - $dE_\nu = j_\nu d\omega dv dt$
- Let us consider thermal emission. Consider a cavity of uniform temperature  $T$  which is a blackbody. This demands
  - $j_\nu^t = \kappa_\nu B_\nu(T)$  – Kirchoff-Planck Law
  - $\kappa_\nu$  is absorption only
  - Strict Thermodynamic Equilibrium applies
  - Thermal absorption and emission independent of angle.

# TE and Stellar Atmospheres

- Energy is transported in a stellar atmosphere
- This means the radiation field is anisotropic
  - There is also a temperature gradient which is demanded by the energy flow (2<sup>nd</sup> Law)
- Strict Thermodynamic Equilibrium cannot hold!
- For convenience we shall assume that local thermodynamic equilibrium holds  $\implies T, N_e$ , etc locally determine occupation numbers.